Rh <sub>II</sub> -2Te <sub>I</sub>	2·76 Å	Te <sub>II</sub> 1Rh <sub>I</sub>	2·64 Å
$-1Te_{II}$	2.61	–2Rh11	2.61
$-2Te_{II}$	2.63	$-4Rh_{II}$	2.63

Table 3.	Structure	factors
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HK0	Fc	$ F_o $	<i>HK</i> 1	Fc	$ F_o $
800	45	41	4.11.1	72	79
3,12,0	-21	30	811	40	40
820	58	58	0,13,1	0	3
840	137	120	1,13,1	36	41
4,12,0	49	60	691	-4	7
780	51	51	831	- 48	46
6,10,0	121	123	771	-15	13
0,14,0	-26	32	2,13,1	- 20	23
1,14,0	5	8	5,11,1	46	46
860	- 49	51	851	36	33
2,14,0	101	111	3,13,1	- 73	80
5,12,0	- 33	33	791	-71	67
920	-30	31	4,13,1	48	48
3,14,0	70	77	911	55	52
940	5	6	6,11,1	8	6
7,10,0	17	17	871	79	71
880	74	74	931	-15	14
4,14,0	-26	30	0,15,1	120	97
6,12,0	- 8	11	1,15,1	- 44	41
960	-27	30	. ,		
0,16,0	-15	15			

–1Rh <sub>I</sub>	2.82
-1Rh1	2.86
-1Rh11	2.91
-2Rh <sub>II</sub>	3.05

The shortest Te–Te distance in the structure is 3.70 Å. In the high temperature (CdI<sub>2</sub>) and low temperature (pyrite) forms of RhTe<sub>2</sub> Geller (1954) reported Rh–6Te =2.65 and 2.64 Å. Geller gives Rh–6Te=2.70 Å and Rh–2Rh=2.83 Å for the high-temperature form (NiAs) of RhTe.

Dr C. Raub made the chemical preparations, Miss Anne Plettinger the single-crystal intensity measurements, and their help is gratefully acknowledged. Part of this work was carried out during a visit to the University of California at San Diego. The writer thanks Professor B. T. Matthias and other La Jolla friends for enjoyable hospitality. The work was in part supported by the Advanced Research Projects Agency.

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# Acta Cryst. (1966). 20, 336 The Crystal Structure of Benzotrifuroxan (Hexanitrosobenzene)\*

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The crystal structure of benzotrifuroxan (hexanitrosobenzene) has been refined by full-matrix leastsquares computations on all positional and thermal parameters to a final weighted R index of 0.040 for 606 reflections of observable intensity. The unit cell is orthorhombic (a = 6.923, b = 19.516, c = 6.518 Å), space group  $Pna2_1$ , with four molecules per unit cell.

The molecule is essentially planar, and the six nitroso substituents have formed three furoxan rings.

### Introduction

There has been considerable interest in the configuration of benzofuroxans. Kaufman & Pickard (1959) cover the classical aspects of this interest in their review article. More recently Harris, Katritzky, Øksne, Bailey & Patterson (1963) have proven the benzofuroxan configuration of this type of compound in solution. Previous X-ray work on this problem consists in two preliminary and somewhat contradictory studies by Hulme (1962) and Gol'der, Todres-Selektor & Bognadov (1961), and work by Britton & Noland (1962) on the structure of 5-chlorobenzofurazan-1-oxide. Britton & Noland found that this compound has the furoxan structure, but their work leaves unsettled some of the important details of the structure. In this paper we report the results of a study whose purpose was to elucidate the structure of benzotrifuroxan in the solid state. In particular, we were concerned with the relative importance of the nitroso and furoxan configurations in this symmetric compound, and with the pertinent bond angles and interatomic distances.

Benzotrifuroxan is, incidentally, a powerful and relatively sensitive explosive.

## Experimental

A sample of benzotrifuroxan was prepared in this Laboratory following the method of Turek (1931). A single crystal suitable for collection of the X-ray intensity data was grown on its mounting fiber by recrystallization from acetone. The longest internal dimension of this crystal was 0.340 mm. The crystal was aligned and preliminary unit-cell dimensions and space group ex-

<sup>\*</sup> This work was performed under the auspices of the U.S. Atomic Energy Commission.

tinctions were determined with precession photographs. The crystal was then transferred to a carefully aligned General Electric Company single-crystal orienter (SCO) equipped with a scintillation counter and a molybdenum X-ray tube.

Cell dimensions were determined from a leastsquares fit of 23  $2\theta$  ( $\lambda$  Mo  $Ka_1=0.70926$  Å) values, greater than 24°, for general *hkl* reflections which were measured on the SCO at 21°C. This least-squares fit gives  $a=6.9234\pm0.0009$ ,  $b=19.5158\pm0.0012$ , c= $6.5180\pm0.0013$  Å for the orthorhombic unit cell. There are four molecules in the unit cell, and the calculated density of 1.901 g.cm<sup>-3</sup> compares reasonably with the measured density of 1.87 g.cm<sup>-3</sup>.

The crystal was mounted with its c axis coincident with the  $\varphi$  axis of the SCO. Background corrections were made by the balanced filter technique. Reflections within a sphere limited by  $2\theta = 50^{\circ}$  were examined by the stationary-counter stationary-crystal technique. Within the quarter of the sphere counted, 1605 spacegroup-allowed reflections were examined. Of the 853 unique reflections in this set, 606 were strong enough to be observed; *i.e.*  $I-Bkgd > 2 \cdot 5 \sqrt{I+Bkgd}$ .

# **Computational details**

Lorentz and polarization corrections were applied. Absorption corrections were made assuming the crystal was box-shaped with sides of lengths 0.145, 0.214, and 0.235 mm (absorption coefficient = 1.87 cm<sup>-1</sup>). The maximum error introduced by this approximation to the true crystal (22 faces) was less than 0.05%. The program described by Larson, Cromer & Roof (1964) was used for these corrections. Equivalent  $F^2$  values were averaged and the observed F was taken as the square root of this average. The average agreement of equivalent reflections was estimated by forming an index

$$\Sigma |F^2 - \overline{F^2}| / \Sigma \overline{F^2}$$
,

where the summation was taken over all reflections which were observed more than once. This index was 0.054. A similar quantity,

$$\Sigma |F - \overline{F^2} \cdot | / \Sigma \overline{F^2} \cdot |$$

based on F, was 0.032.

All least-squares calculations were made with the use of the full matrix. The least-squares parameters included, in addition to atomic position and thermal parameters, one scale factor and one parameter to allow for secondary extinction. The secondary extinction parameter was found to be negligible and was set equal to zero in the later stages of the refinement. The quantity minimized was  $\Sigma w(F_o - F_c^*)^2$ , where

$$|F_c^*| = |F_c|/k \left\{ 1 + g \left[ \frac{2(1 + \cos^4 2\theta)}{(1 + \cos^2 2\theta)^2} \right] Lp |F_c|^2 \right\}^{\pm}$$

in which k = scale factor, g = extinction parameter (Zachariasen, 1963), Lp = Lorentz and polarization factors, and  $F_c$  is the calculated structure factor, and where  $w = w_E/(|F_o| + 0.02|F_o|^2)$  for all observed reflections and w = 0.0 for unobserved reflections. The term  $w_E$  is the weight based on counting statistics as described by Evans (1961). Anisotropic thermal parameters were in the form  $\exp[-(h^2B_{11}+k^2B_{22}+l^2B_{33}+hkB_{12}+hlB_{13}+klB_{23})]$ . The *R* index is defined as  $(\Sigma w||F_o| - |F_c^*||/\Sigma w|F_o|)$ .

The estimated standard deviations were calculated from

$$\sigma_j = \sqrt{a^{jj} \left[ \frac{\Sigma w(|F_o| - |F_c^*|)^2}{m - n} \right]}$$

where *m* is the number of observed reflections, *n* is the number of parameters, and  $a^{jj}$  is the *jj* element of the inverse matrix. The atomic scattering factors used are those given in *International Tables for X-Ray Crystallography* (1962).

All calculations were performed on IBM 7094 or 7030 computers, with codes written by the authors unless otherwise indicated.

### Determination and refinement of the structure

The space group is not uniquely determined by systematic extinctions (0kl with (k+l) odd, hol with h odd)and could be either  $Pna2_1$  or Pnam. A three-dimensional sharpened Patterson map was computed. Inspection of the Patterson map and statistical tests of the data (Howells, Phillips & Rogers, 1950; Ramachandran & Srinivasan, 1959) indicated the proper space group to be  $Pna2_1$ . The orientation of a molecule could be derived from the Patterson map; however, the x and y coordinates of atoms were not obvious.

Since the general position set for space group  $Pna2_1$ is fourfold and because so much was known about the geometry of the molecule, the problem seemed ideally suited to a vector verification method for determining a trial structure. A code was written which would start with a fixed array of points and then rotate and translate this array into any of its possible orientations in the unit cell. The symmetry related points were generated and the value of the sharpened Patterson map at all interpoint vectors was determined. Scaling of the values for the assumed atomic numbers of the atoms in the vectors, and selection of the lowest value determined from an entire set of vectors, yielded a number indicating the consistency of the model with the Patterson map. This procedure amounts to calculating the value of the minimum function (Buerger, 1959) at each 'atom site' and then selecting the lowest of these values as the index of correctness of the postulated structure.

A search was conducted for likely locations for a benzene group. After many possible benzene group positions were found, nitrogen atoms were added at 1.4 Å from the carbon atoms and several possible trial structures were obtained.

Close inspection of these trial structures revealed that they were all the same except for choice of origin or polar direction in the unit cell.

One set of carbon and nitrogen positions was used to calculate an observed Fourier map. The extra peaks on this map determined the trial positions for the oxygen atoms in the least-squares refinement. With the assumption of isotropic atoms the problem converged rapidly to an R index of 0.10. All atoms were then allowed to become anisotropic, and the refinement continued until all  $\Delta \xi_i / \sigma(\xi_i)$  were  $< 1.9 \times 10^{-2}$  for all parameters  $\xi$ . The final value of R was 0.040.

The final least-squares parameters and their standard deviations are given in Table 1 and the corresponding values of  $10 k |F_o|$ ,  $10 k |F_c^*|$ , and  $\alpha$  are given in Table 2. The standard deviation of the electron density ( $\sigma_{\rho}$ ) in the structure is estimated to be  $4.73 \times 10^{-2}$ e.Å<sup>-3</sup> by Cruickshank's (1949) formula. Observed and difference Fourier maps of the final structure were calculated. The peaks and valleys of the difference map indicated no significant errors in the structure.

# Detailed description of the structure

Interatomic distances before correction for thermal motion are given in Table 3. The closest intermolecular distance is from N(1) to O(5') of the molecule at z+1,

which is  $2.997 \pm 0.006$  Å. Intramolecular bond angles are given in Fig. 1 (standard deviations for these angles range from 0.4 to 0.6 degrees). Standard deviations for



Table 1. The final least-squares parameters and their estimated standard deviations for benzotrifuroxane

The standard deviations apply to the rightmost digits of the least-squares parameters. The anisotropic temperature factor is of the form:  $\exp \left[-(B_{11}h^2 + B_{22}k^2 + B_{33}l^2 + B_{12}hk + B_{13}hl + B_{23}kl)\right]$ ; and B is  $\frac{4}{3}(B_{11}a^2 + B_{22}b^2 + B_{33}c^2)$  for an anisotropic atom (Hamilton, 1959).

		5.		55-7			,			
Atom	x	у.	Ζ	$B_{11} \times 10^4$	$B_{22} \times 10^5$	$B_{33} \times 10^4$	$B_{12} \times 10^{5}$	$B_{13} \times 10^{4}$	$B_{23} \times 10^4$	$B(\text{\AA})^2$
<b>C</b> (1)	$0.4165 \pm 8$	$0.0995 \pm 2$	0	$217 \pm 13$	<b>229 ±</b> 16	$245 \pm 17$	$-250 \pm 82$	$-1 \pm 28$	$-10 \pm 9$	3.94
C(2)	0.6009 <u>+</u> 8	0·0743 ± 2	$-0.0691 \pm 13$	$186 \pm 12$	$223 \pm 15$	$282 \pm 17$	51 ± 77	$-77 \pm 27$	$-8 \pm 10$	3.92
C(3)	$0.6832 \pm 7$	0·0950±3	$-0.2554 \pm 14$	$175 \pm 12$	$235 \pm 16$	$301 \pm 18$	$-13 \pm 79$	$12 \pm 28$	$-24 \pm 11$	4.02
C(4)	0·5782 <u>+</u> 7	0·1413±3	$-0.3817 \pm 12$	$201 \pm 14$	249 ± 16	$221 \pm 17$	$-198 \pm 79$	$64 \pm 26$	$-18 \pm 9$	3.80
C(5)	0·3975±7	0·1693 ± 2	$-0.3187 \pm 13$	$207 \pm 13$	210±16	$236 \pm 17$	$-40 \pm 75$	$-15 \pm 27$	$-8 \pm 9$	3.73
C(6)	$0.3216 \pm 7$	$0.1475 \pm 2$	$-0.1255 \pm 13$	$171 \pm 12$	229 ± 16	$272 \pm 17$	21 ± 78	5 ± 27	$-36 \pm 9$	3.79
N(1)	0.3219+6	0.0851 + 2	0.1677 + 11	223 + 11	343 + 17	275 + 15	-83+73	40 + 26	12+10	4.72
N(2)	$0.7216 \pm 7$	$0.0317 \pm 2$	0.0216 + 12	210 + 12	275 + 16	$326 \pm 18$	$-34 \pm 71$	$-29 \pm 26$	$6 \pm 9$	4.58
N(3)	$0.8558 \pm 6$	$0.0676 \pm 3$	$-0.2843 \pm 13$	$218 \pm 12$	$335 \pm 17$	$363 \pm 19$	$6 \pm 68$	$57 \pm 29$	$-34\pm11$	5.15
N(4)	$0.6153 \pm 7$	$0.1657 \pm 2$	$-0.5680 \pm 12$	$267 \pm 13$	$245 \pm 14$	$260 \pm 16$	$-262\pm76$	$26 \pm 27$	$-16\pm 8$	4.42
N(5)	0·3245±7	$0.2106 \pm 2$	$-0.4545 \pm 12$	$276 \pm 13$	$311 \pm 16$	$267 \pm 16$	$207 \pm 80$	$-13 \pm 27$	13±9	4·85
N(6)	0·1588 ± 7	$0.1653 \pm 2$	$-0.0291 \pm 13$	$199 \pm 11$	$339 \pm 16$	$301 \pm 16$	$-7 \pm 76$	$20 \pm 24$	$-20\pm10$	4·70
<b>O</b> (1)	0·1548±5	$0.1252 \pm 2$	$0.1610 \pm 12$	$214 \pm 10$	409 ± 15	$296 \pm 13$	$28 \pm 61$	82 ± 20	$4\pm 8$	5.12
O(2)	$0.7192 \pm 6$	$-0.0014\pm 2$	$0.1787 \pm 13$	$277 \pm 11$	$389 \pm 15$	$388 \pm 14$	$180 \pm 68$	$-61 \pm 25$	89±9	5.95
O(3)	$0.8896 \pm 5$	$0.0274 \pm 2$	$-0.1137 \pm 12$	$208 \pm 9$	$330 \pm 12$	$379 \pm 15$	$224 \pm 57$	$-28 \pm 25$	$-6 \pm 9$	5.15
O(4)	0·7412±6	0·1591 ± 2	$-0.6934 \pm 12$	$298 \pm 10$	$402 \pm 15$	$277 \pm 14$	$-344 \pm 67$	$206 \pm 23$	- 34 <u>+</u> 8	5.51
O(5)	0·4533±6	0·2122 ± 2	$-0.6154 \pm 11$	$301 \pm 11$	$330 \pm 13$	279 <u>+</u> 13	$104 \pm 63$	$1 \pm 24$	9±8	5.18
O(6)	$0.0262 \pm 6$	$0.2035 \pm 2$	$-0.0656 \pm 12$	$247 \pm 10$	$405 \pm 14$	$437 \pm 17$	722 ± 64	$15 \pm 23$	$18 \pm 10$	6.11

these distances and bond angles were computed using the entire variance—covariance matrix. Examination of Table 3 and Figs.1 and 2 clearly indicates the benzotrifuroxan structure as opposed to any other of the possible configurations.

The atomic positional parameters were transformed into an isometric set and the plane of the molecule was determined to be 0.46634X + 0.76605Y + 0.44237Z - 2.86851 = 0. The distances of the atoms from this plane are given in Table 4.

The anisotropic thermal parameters were transformed to obtain the axes of the ellipsoids of thermal vibration and the directions of these axes with respect to the crystallographic axes. These results are presented in Table 5. The standard deviations were computed using the entire variance-covariance matrix. Rigid body calculations were carried out for the following groups of atoms: (1) the benzene ring, (2) all of the carbon and nitrogen atoms, and (3) the benzotrifurazan ring system. The method of Cruickshank (1956a) was employed using a code provided by Trueblood (1963). The center of mass of the rigid body was taken as the unweighted average of the atomic positions. Oxygen atoms (2), (4), and (6) were clearly eliminated from the rigid body by examination of their ellipsoids of thermal vibration. Examination of the thermal and oscillatory motions from the three rigid body calculations showed that the benzene ring was the only group which had reasonable rigid body motions.

The rigid body calculations are strictly valid only when the whole molecule may be regarded as a rigid

# Table 2. Observed and calculated magnitudes of the structure factors for benzotrifuroxan

The column headings are k,  $10k|F_0|$ ,  $10kF_c$  (see text), and  $\alpha$ . A minus sign preceding  $10k|F_0|$  means 'less than'.

body, or when the internal vibrations are small enough to be neglected. This molecule does not meet these requirements. Likewise the assumption that all bonded atoms are vibrating in phase seems somewhat unlikely, especially for the atoms in the furazan rings. However, in the absence of facilities for a more detailed and rigorous treatment, a combination of these approximations was tried. Bond distance corrections were calculated for the carbon skeleton by the method of Cruickshank (1956b, 1961) and for the rest of the atoms assuming adjacent atoms vibrate in phase (Busing & Levy, 1964). Support for this assumption comes from the fact that the thermal vibration of all atoms increases perpendicular to the bonds as the distance from the center of the molecule increases. The corrected bond distances are given in Table 3, and a diagram of the molecule showing these distances is shown in Fig.2.

Bond number calculations were made (Pauling, 1960) using the average corrected bond distances from Table 3. These calculations indicate bond numbers of 1.27 for the C-C, 1.46 for the type C(1)-N(1), 1.34 for the type C(2)-N(2), 1.74 for the type N(2)-O(2), 0.82 for the type N(2)-O(3), and 1.03 for the type N(3)-O(3) bonds.

1 abic 5. Interatornic distances in Denzotrijaroza	Та	able 3. Interat	tomic distances	: in l	benzotrifuroxan	!
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Atoms	Length (Å) Least-squares	Length (Å) Corrected for thermal motion
C(1) - C(2)	$1.440 \pm 0.007$	1.454
C(2) - C(3)	$1.401 \pm 0.008$	1.416
C(3) - C(4)	$1.422 \pm 0.007$	1.435
C(4) - C(5)	$1.425 \pm 0.007$	1.439
C(5) - C(6)	$1.429 \pm 0.008$	1.444
C(6) - C(1)	$1.407 \pm 0.007$	1.419
C(1) - N(1)	$1.305 \pm 0.007$	1.316
C(3) - N(3)	$1.322 \pm 0.006$	1.335
C(5) - N(5)	$1.300 \pm 0.007$	1.316
C(2) - N(2)	$1.318 \pm 0.007$	1.324
C(4) - N(4)	$1.330 \pm 0.007$	1.336
C(6) - N(6)	$1.336 \pm 0.007$	1.346
N(1) - O(1)	$1.398 \pm 0.005$	1.403
N(3) - O(3)	$1.380 \pm 0.007$	1.384
N(5)-O(5)	$1.377 \pm 0.005$	1.380
N(2) - O(2)	1·211 ± 0·007	1.236
N(4)-O(4)	$1.201 \pm 0.006$	1.225
N(6)-O(6)	$1.207 \pm 0.005$	1.235
N(2) - O(3)	$1.462 \pm 0.006$	1.469
N(4)-O(5)	$1.475 \pm 0.006$	1.486
N(6) - O(1)	$1.466 \pm 0.007$	1.473

# Table 4. Distances of atoms from plane of molecule

Atom	Distance (A)	Atom	Distance (Å)	
<b>C</b> (1)	-0.036	N(4)	-0.042	
C(2)	-0.016	N(5)	0.012	
C(3)	0.021	N(6)	0.031	
C(4)	0.010	O(1)	-0.033	
C(5)	0.027	O(2)	-0.02	
<b>C</b> (6)	0.013	O(3)	0.086	
N(1)	-0.074	O(4)	-0.096	
N(2)	-0.002	O(5)	0.007	
N(3)	0.082	O(6)	0.020	

# Table 5. Magnitudes and direction angles, relative to the<br/>crystallographic axes, of the principal axes<br/>of the vibration ellipsoids

Atom	Axis i	$B_i$ (Å <sup>2</sup> )	a (°)	b (°)	<b>c</b> (°)
<b>C</b> (1)	1 2 3	$4.6 \pm 0.3$ $3.0 \pm 0.2$ $4.2 \pm 0.3$	$38 \pm 18 \\ 60 \pm 8 \\ 111 \pm 26$	$123 \pm 8$ $33 \pm 8$ $90 \pm 19$	$72 \pm 30$ $79 \pm 11$ $21 \pm 26$
C(2)	1 2 3	$\begin{array}{c} 3 \cdot 2 \pm 0 \cdot 2 \\ 3 \cdot 4 \pm 0 \cdot 2 \\ 5 \cdot 1 \pm 0 \cdot 3 \end{array}$	$28 \pm 39 \\ 77 \pm 74 \\ 114 \pm 7$	$107 \pm 81$ 19 \pm 73 98 \pm 8	$69 \pm 23 \\ 76 \pm 30 \\ 26 \pm 7$
C(3)	1 2 3	$\begin{array}{c} 3 \cdot 3 \pm 0 \cdot 2 \\ 3 \cdot 3 \pm 0 \cdot 2 \\ 5 \cdot 3 \pm 0 \cdot 3 \end{array}$	$3 \pm 30 \\ 90 \pm 392 \\ 87 \pm 7$	$\begin{array}{c} 88 \pm 371 \\ 19 \pm 27 \\ 109 \pm 7 \end{array}$	93±127 71±22 19±7
C(4)	1 2 3	$\begin{array}{c} 4 \cdot 8 \pm 0 \cdot 3 \\ 3 \cdot 3 \pm 0 \cdot 3 \\ 3 \cdot 2 \pm 0 \cdot 3 \end{array}$	$52 \pm 8$ 77 ± 72 139 ± 33	$124 \pm 8$ $36 \pm 35$ $102 \pm 77$	$56 \pm 8$ 57 ± 70 52 ± 66
C(5)	1 2 3	$3 \cdot 9 \pm 0 \cdot 3$ $3 \cdot 1 \pm 0 \cdot 2$ $4 \cdot 1 \pm 0 \cdot 3$	$36 \pm 53$ $81 \pm 14$ $124 \pm 54$	$106 \pm 16 \\ 17 \pm 15 \\ 96 \pm 20$	$58 \pm 53 \\ 76 \pm 14 \\ 35 \pm 50$
<b>C</b> (6)	1 2 3	$\begin{array}{c} 3 \cdot 3 \pm 0 \cdot 2 \\ 3 \cdot 0 \pm 0 \cdot 2 \\ 5 \cdot 1 \pm 0 \cdot 3 \end{array}$	$13 \pm 34 \\ 103 \pm 34 \\ 90 \pm 8$	$\begin{array}{c} 79 \pm 30 \\ 32 \pm 14 \\ 119 \pm 6 \end{array}$	$84 \pm 18$ $61 \pm 7$ $29 \pm 6$
N(1)	1 2 3	$4.0 \pm 0.3$ $5.4 \pm 0.3$ $4.8 \pm 0.3$	$36 \pm 13$ $94 \pm 15$ $54 \pm 13$	$74 \pm 9 \\ 23 \pm 18 \\ 106 \pm 22$	$122 \pm 12$ $68 \pm 20$ $40 \pm 18$
N(2)	1 2 3	$4.0 \pm 0.2$ $4.2 \pm 0.3$ $5.6 \pm 0.3$	$18 \pm 39$ $105 \pm 46$ $100 \pm 8$	$74 \pm 46 \\ 18 \pm 43 \\ 83 \pm 9$	$83 \pm 12$ $100 \pm 11$ $12 \pm 9$
N(3)	1 2 3	$4.0 \pm 0.3$ $4.7 \pm 0.3$ $6.7 \pm 0.4$	$24 \pm 14 \\ 68 \pm 15 \\ 80 \pm 6$	$105 \pm 13$ $32 \pm 10$ $117 \pm 7$	$109 \pm 8 \\ 68 \pm 8 \\ 29 \pm 6$
N(4)	1 2 3	$5.5 \pm 0.3 \\ 3.4 \pm 0.2 \\ 4.4 \pm 0.3$	$31 \pm 9$ $71 \pm 6$ $114 \pm 12$	$114 \pm 5$ 25 \pm 6 97 ± 11	71 ± 11 74 ± 11 25 ± 12
N(5)	1 2 3	$5.6 \pm 0.3 \\ 4.1 \pm 0.3 \\ 4.8 \pm 0.3$	$33 \pm 11$ $112 \pm 10$ $113 \pm 16$	$57 \pm 10$ $47 \pm 12$ $60 \pm 16$	$85 \pm 14$ 129 ± 16 40 ± 16
N(6)	1 2 3	$\begin{array}{c} 3 \cdot 8 \pm 0 \cdot 2 \\ 5 \cdot 6 \pm 0 \cdot 3 \\ 4 \cdot 6 \pm 0 \cdot 3 \end{array}$	$9 \pm 13$ 94 $\pm 7$ 82 $\pm 14$	$92 \pm 11$ $44 \pm 11$ $46 \pm 11$	$99 \pm 11$ $134 \pm 11$ $45 \pm 12$
<b>O</b> (1)	1 2 3	$\begin{array}{c} 3 \cdot 7 \pm 0 \cdot 2 \\ 6 \cdot 2 \pm 0 \cdot 2 \\ 5 \cdot 4 \pm 0 \cdot 2 \end{array}$	$29 \pm 5$ $85 \pm 7$ $61 \pm 6$	$90 \pm 4$ $9 \pm 13$ $99 \pm 13$	$119 \pm 5$ $82 \pm 12$ $30 \pm 7$
O(2)	1 2 3	$5.6 \pm 0.2 \\ 3.6 \pm 0.2 \\ 8.6 \pm 0.3$	$24 \pm 6$ $114 \pm 5$ $92 \pm 4$	$71 \pm 5 \\ 47 \pm 3 \\ 50 \pm 2$	$104 \pm 4$ $127 \pm 3$ $40 \pm 2$
O(3)	1 2 3	$\begin{array}{c} 3.7 \pm 0.2 \\ 5.3 \pm 0.2 \\ 6.5 \pm 0.3 \end{array}$	$25 \pm 5$ $67 \pm 6$ $98 \pm 6$	$114 \pm 5$ 26 \pm 6 99 \pm 10	86±5 79±10 12±10
O(4)	1 2 3	$\begin{array}{c} 8 \cdot 0 \pm 0 \cdot 3 \\ 5 \cdot 3 \pm 0 \cdot 2 \\ 3 \cdot 3 \pm 0 \cdot 2 \end{array}$	$50 \pm 3$ $62 \pm 5$ $126 \pm 3$	$124 \pm 4$ $34 \pm 4$ $87 \pm 4$	59±3 73±4 36±3
O(5)	1 2 3	$5.9 \pm 0.2 \\ 5.1 \pm 0.2 \\ 4.6 \pm 0.2$	$20 \pm 12$ $109 \pm 13$ $83 \pm 11$	$\begin{array}{c} 70 \pm 11 \\ 39 \pm 20 \\ 122 \pm 20 \end{array}$	$86 \pm 12 \\ 57 \pm 22 \\ 33 \pm 21$
O(6)	1 2 3	$3.4 \pm 0.2$ $7.9 \pm 0.3$ $7.0 \pm 0.3$	$35 \pm 2$ 66 \pm 6 114 \pm 6	$125 \pm 2$ $51 \pm 8$ $122 \pm 9$	$88 \pm 3$ $48 \pm 12$ $42 \pm 12$

These bond numbers indicate a substantial amount of resonance in the ring system. Weakening of the C–C bonds in the benzene ring by this resonance was found in benzotrifuroxan just as it was in 1,3,5-triamino-2,4,6-trinitrobenzene (Cady & Larson, 1965).

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# The Structure of Fibrous Sulphur

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The structure of fibrous sulphur has been determined by a systematic use of the Fourier transform of its helical molecules. From this analysis, the helical molecules appeared to have two *incommensurable* repetition lengths in the direction of the helical axis, one being due to the succession of the atoms in that direction, the other to the pitch of the helix. A trial to substitute the two repetition lengths by a single one would result in an infinite repetition length. For this reason the structure is not crystallographic in the classical sense of the word, so we have chosen an improper 'unit cell' of orthorhombic symmetry, containing four helical molecules of supposedly infinite length, with the dimensions  $a=8\cdot11$  Å and  $b=9\cdot20$  Å, c= indeterminate. The directions of a and b are normal to the helical axis, but only b is properly a repetition length.

In the direction of a the structure consists of alternating layers of right and left handed helices. Within each layer the helical molecules are screwed in or out, with respect to their neighbours, over half the atomic period along the molecule. This explains a peculiar extinction rule.

The molecular parameters, *i.e.* the bond distance, the bond angle and the dihedral angle, are compared with those of other sulphur molecules.

### Introduction

Several physical properties of liquid sulphur suggest that above 160 °C it contains long chain molecules (see *e.g.* Schenk, 1956). One of the strongest arguments is provided by the X-ray fibre diagram, obtained from fibres made by highly stretching amorphous sulphur quenched from above 300 °C. This diffraction pattern was discovered by Trillat & Forestier (1932); after being wrongly interpreted initially by Meyer & Go (1934) as a onecomponent pattern, it was later proved (Prins, Schenk

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& Hospel, 1956) to be due to two components of which only one is soluble in carbon disulphide. The reflexions due to the soluble part were shown (Prins, Schenk & Wachters, 1957) to belong to a metastable ring modification  $S_y$  (De Haan, 1958). The reflexions of the carbon-disulphide-insoluble part evidently formed a fibre diagram. The corresponding metastable modification was called  $S_y$ .

Schenk (1956), comparing  $S_{\psi}$  with hexagonal selenium, proposed a preliminary unit cell containing 10 atoms placed in three turns of a cylindrical-helical mol-